

# SIMPLE MODEL FOR THE CALCULATION OF THE COEFFICIENT OF SELF-DIFFUSION IN A LIQUID

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The coefficient of self-diffusion in three-dimensional classical liquid is computed approximately from the hierarchy of kinetic equations for the time-correlation functions (TCF).

The coefficient of self-diffusion  $D_s$  of a molecule can be obtained from the TCF  $\pi(t) = \langle \mathbf{p}_1(0) \mathbf{p}_1(t) \rangle / \langle \mathbf{p}_1^2 \rangle$  of momentum  $\mathbf{p}_1(t)$  at time  $t$  of a chosen particle in a fluid by  $D_s = (kT/m) \lim_{z \rightarrow +0} \tilde{\pi}(z)$ , where  $\tilde{\pi}(z) = \int_0^\infty dt \exp(-zt) \pi(t)$ , where  $m$  is the mass of a particle and  $kT$  the thermal energy. Here we shall obtain the formula for  $D_s$  from the small- $z$  behavior of the hierarchy of the equations for TCF.

Zwanzig was the first to succeed in deriving a kinetic equation from the Liouville equation [1]:  $\pi'(t) = - \int_0^t d\tau K(\tau) \pi(t - \tau)$ , where the memory function is

$$K(\tau) = \langle \mathbf{p}_1 \hat{\mathcal{L}} \exp\{i\mathcal{L}_{22}^{(1)}\tau\} \hat{\mathcal{L}} \mathbf{p}_1 \rangle / \langle \mathbf{p}_1^2 \rangle, \quad \mathcal{L}_{22}^{(1)} = P \hat{\mathcal{L}} P, \quad \hat{\mathcal{L}} = -i\hat{L}, \quad \hat{\mathcal{L}} = \sum_{j=1}^N \frac{1}{m} \mathbf{p}_j \nabla_j - \sum_{i \neq j=1}^N \nabla_j u(i, j) \nabla_{\mathbf{p}_j},$$

$$P = 1 - \Pi, \quad \Pi = \frac{\mathbf{p}_1 \langle \mathbf{p}_1 \rangle}{\langle \mathbf{p}_1^2 \rangle},$$

$\Pi$  is the projection operator,  $\hat{L}$  is the ordinary Liouville operator and  $u(i, j)$  is the interparticle potential. If we use the identity [2]

$$\exp\{t(A+B)\} = \exp(tA) + \int_0^t du \exp\{(t-u)A\} B \exp\{u(A+B)\},$$

for the arbitrary operators  $A$  and  $B$ , we can obtain an expansion of  $K(t)$ :

$$K(t) = \omega^2 f(t) + \sum_{n=1}^{\infty} (-1)^{n+1} \int_0^t dt_1 \int_0^{t_1} dt_2 \dots \int_0^{t_{n-1}} dt_n \pi'(t-t_1) \pi'(t_1-t_2) \dots \pi'(t_{n-1}-t_n) \pi''(t_n),$$

where  $\omega^2 = \langle F_1^2 \rangle / \langle \mathbf{p}_1^2 \rangle$ ,  $f(t) = \langle F_1 \exp(i\hat{\mathcal{L}}t) F_1 \rangle / \langle F_1^2 \rangle$ ,  $F_1 = -\sum_{j>1}^N \nabla_1 u(1, j)$  is the total force on a chosen particle.

Now one can write the Laplace transform of  $K(t)$  in the form

$$\tilde{K}(z) = \omega^2 \tilde{f}(z) + z \sum_{n=1}^{\infty} \{1 - z \tilde{\pi}(z)\}^{n+1}. \quad (1)$$

Then for  $z \rightarrow +0$  we obtain  $\tilde{K}(z) \approx \omega^2 \tilde{f}(z)$  and  $\tilde{\pi}(z) \approx \{z + \omega^2 \tilde{f}(z)\}^{-1}$ . Following Zwanzig [1] we may write the kinetic equation for the TCF  $f(t)$ :  $f'(t) = - \int_0^t d\tau V(\tau) f(t - \tau)$ , where

$$V(\tau) = \langle F_1 \hat{\mathcal{L}} \exp\{i\mathcal{L}_{22}^{(2)}\tau\} \hat{\mathcal{L}} F_1 \rangle / \langle F_1^2 \rangle, \quad \mathcal{L}_{22}^{(2)} = Q \hat{\mathcal{L}} Q, \quad Q = 1 - R, \quad R = \frac{F_1 \langle F_1 \rangle}{\langle F_1^2 \rangle},$$